

Non-equilibrium Plasma in Multibubble Sonoluminescence



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Acoustic cavitation, Göttingen and me



01/2008: I was appointed at CEA / ICSM to work on sonoluminescence combined to sonochemistry

But: building under construction, no lab

Collaboration with MPI-KG Potsdam: there, SBSL, at ICSM, MBSL



Looking at the literature: specialists in SBSL nearby, in Göttingen

03/2008: Visit on a snowy day

T. Kurz explained me everything you should know to get started

Prof. Lauterborn was not there but I could admire some of his paintings



Akustische Blasenfallen
für die
Sonolumineszenz



Diplomarbeit
vorgelegt von
Reinhard Geisler
aus
Göttingen

angefertigt im
Dritten Physikalischen Institut
der Georg-August-Universität zu Göttingen

1998

06/2008: my first ESS (ESS-16)

Robert Mettin going from a poster to the next one and giving advice to every student

Sharing knowledge as common trait in their group?

10/2011: Workshop Kavitation Drübeck

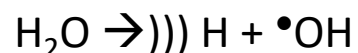
Finally I met Prof. Lauterborn!



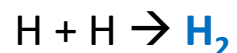
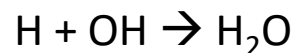
12/2013: Göttingen, experiments & visit of the labs

The very violent collapse of cavitation bubbles leads to the formation of a plasma

Radical generation in the sonochemical plasma:

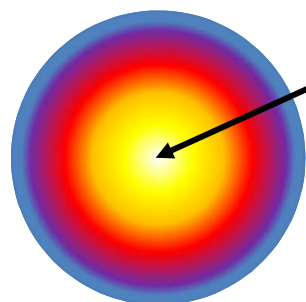


Subsequent recombination :



} Basis of sonochemical activity in aqueous solutions

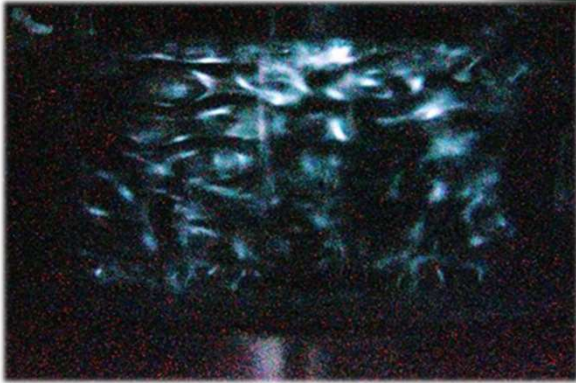
3 zones of sonochemical activity:



Hot bubble core, plasma (thousands of degrees, hundreds of bars) (< 1 μm , 100 ps)

Overheated bubble-solution interface (hundreds of degrees)

Volume of the solution

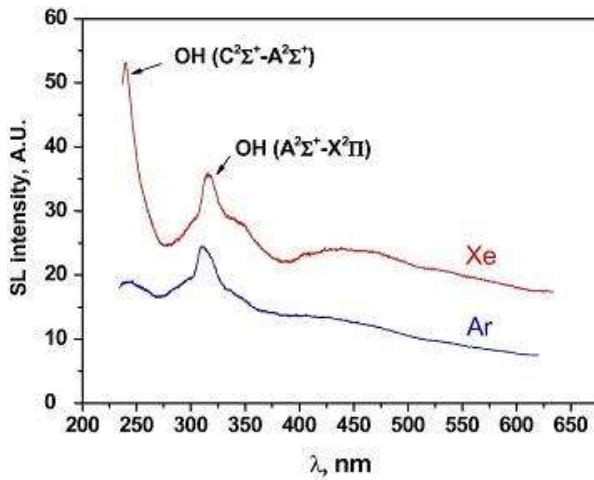
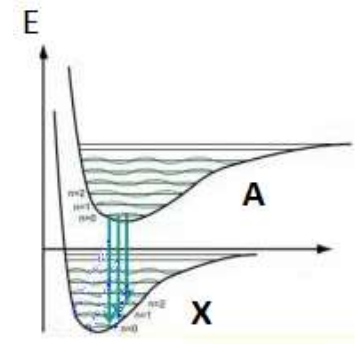


Light emitted by the sonochemical plasma
probe of the conditions reached inside bubbles

Intense continuum + emissions from electronically excited species

- Hints on formed species, on mechanisms of reactions
- Emissions are characteristic of the species and its excitation state

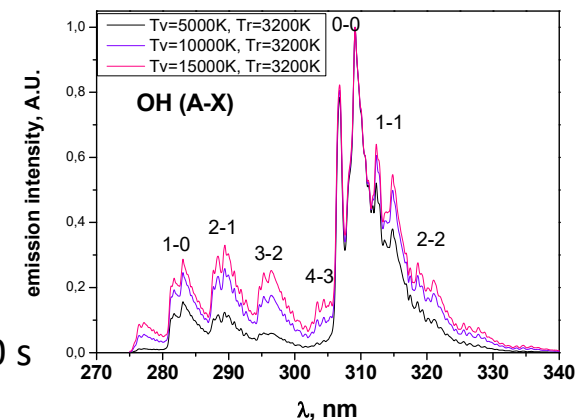
Estimation of temperatures



Average on time and space

Wavelength in water: 4.1 mm
at 362 kHz

Typical acquisition time of a spectrum: 10-600 s
Plasma lifetime < 1 μs
US period: 50 μs at 20 kHz, 2.8 μs at 362 kHz

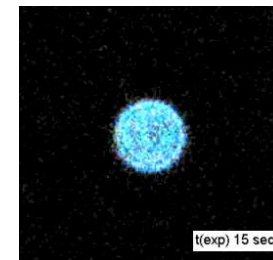


- 1933: Marinesco and Trillat observed that a photographic plate submitted to US went blurred
- 1934: Frenzel and Schultes noticed a faint luminescence, responsible for the blurring
- Many measurements (PMT...) : cyclic flashes of light, frequency-dependent etc. But averages
- Revolution in 1962 (Yosioka and Omura) - 1970 (Temple) – 1990 (Gaitan and Crum) : discovery of **single-bubble sonoluminescence (SBSL)**



Crum, 1994

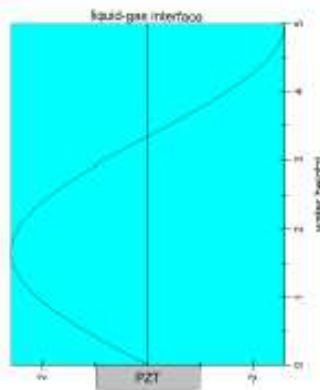
SBSL made it possible to measure bubble dynamics, light emission etc.



How to observe a single bubble



- Avoid 'natural' cavitation:
→ degas solution to be **under the cavitation threshold**
- Introduce and stabilise a bubble
→ bubble trap: **standing wave**



$$h = \frac{(2n+1)\lambda}{4}$$

$$\lambda = \frac{c}{f}$$

c sound speed, in water at 20°C $c=1480 \text{ m}\cdot\text{s}^{-1}$

f US frequency

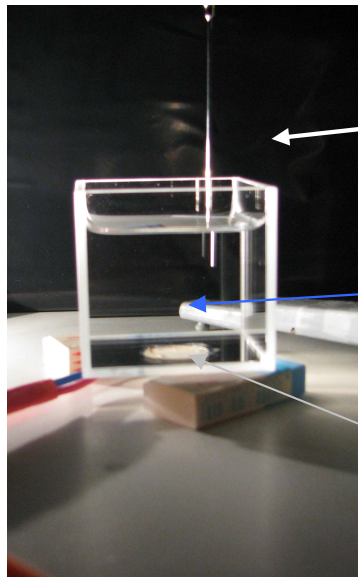
At 20 kHz : $\lambda=7.4 \text{ cm}$

At 300 kHz : $\lambda=4.9 \text{ mm}$

By controlling the amount of dissolved gas and the acoustic pressure, it is possible to stabilise a single bubble for hours.

PhD Julia Schneider (2008-2012)

First SBSL reactor:
inspired by Göttingen

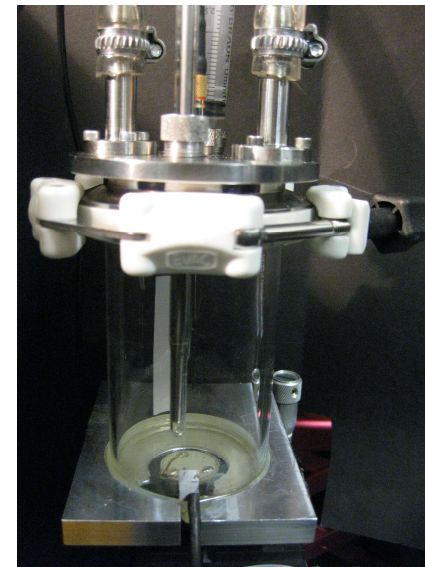


Hydrophone

Single bubble
trapped in the
stationary wave

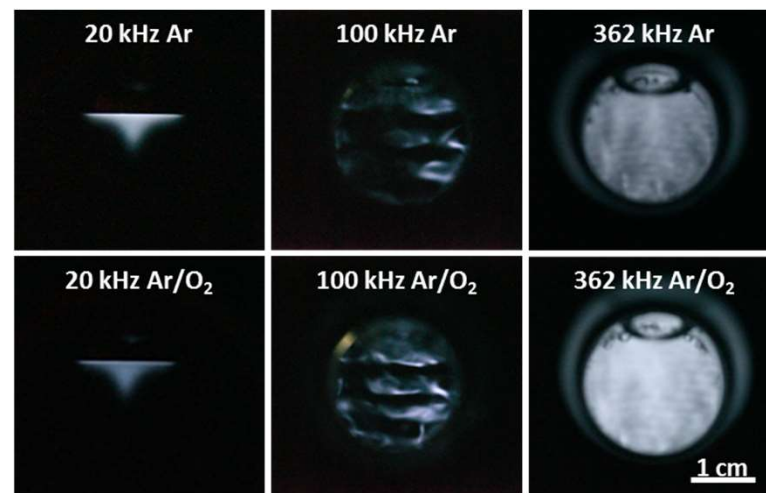
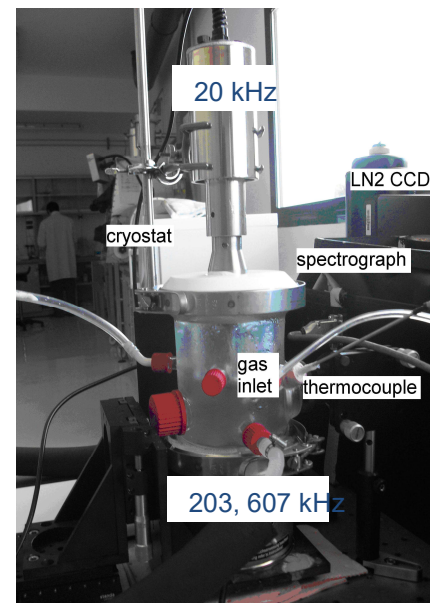
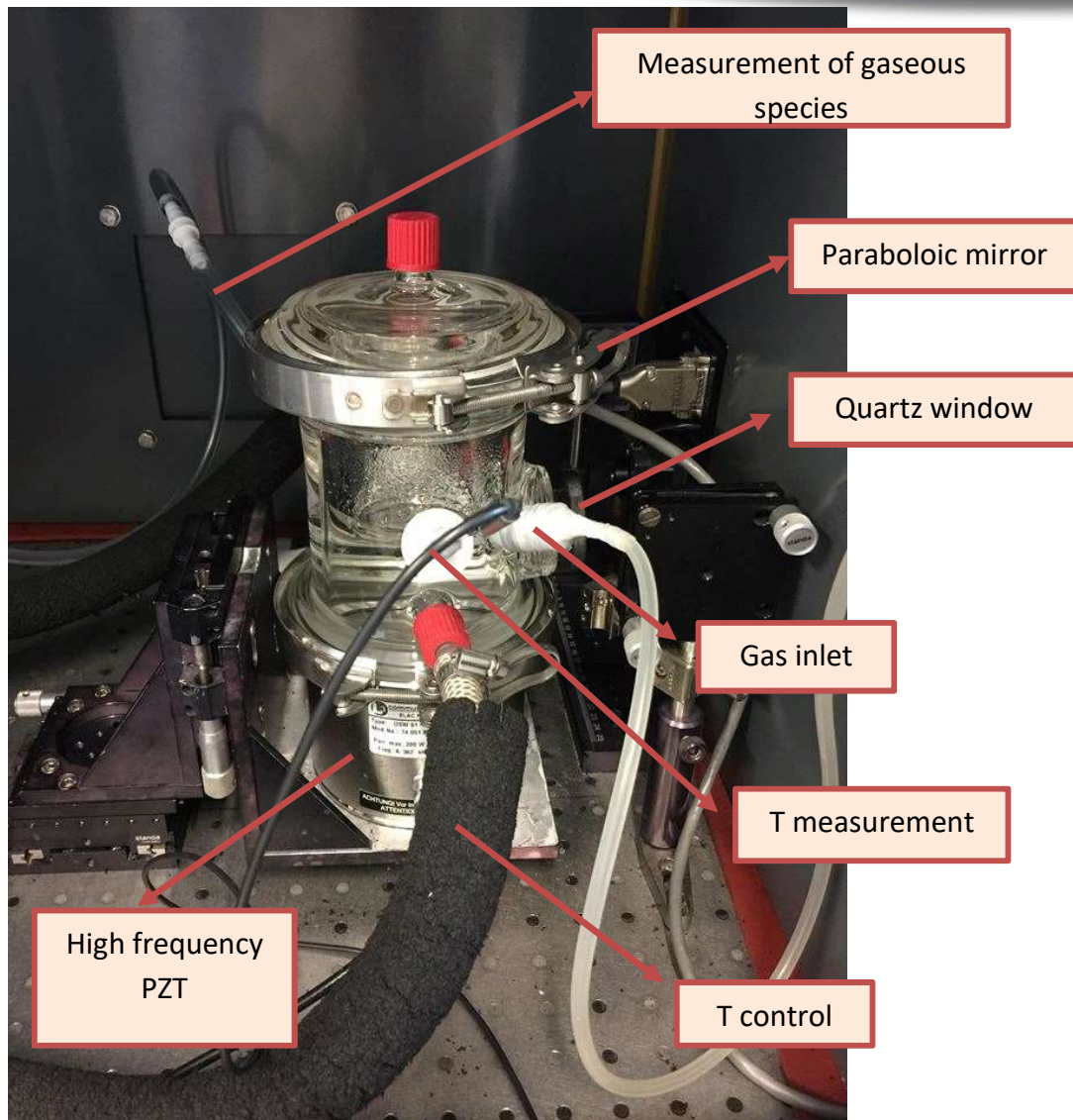
Piezo-element

SBSL in controlled conditions:
temperature, gas nature et partial
pressure, acoustic pressure
106 mL; 27.6 kHz



*SBSL in water (air)
(10 s exposure time)*

MBSL set-up, Marcoule



250 mL, continuous gas flow, 10 °C
Frequency: 20 / 100 / 200 / 362 / 600 / 1057 kHz

Plasma?



Camera Lucida, common work of artists and scientists,
Göttingen (H_2SO_4 , Xe)

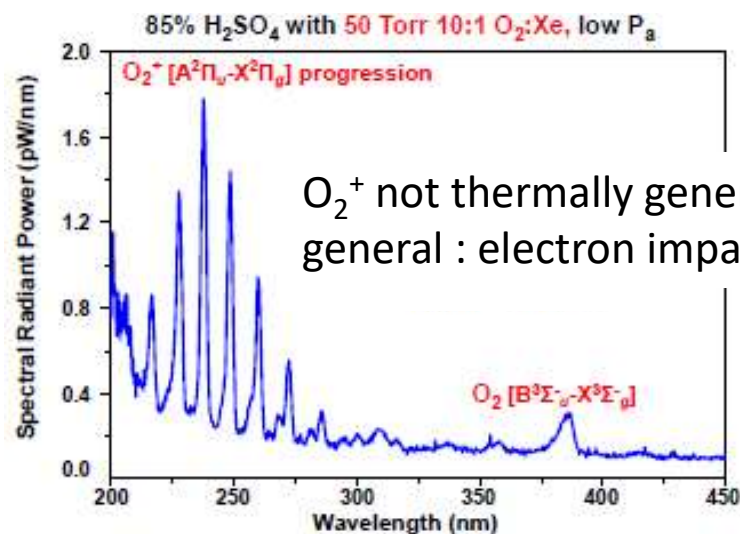
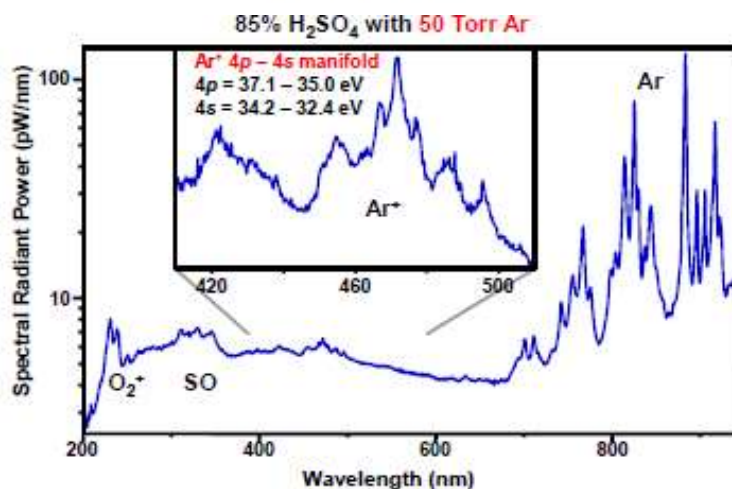
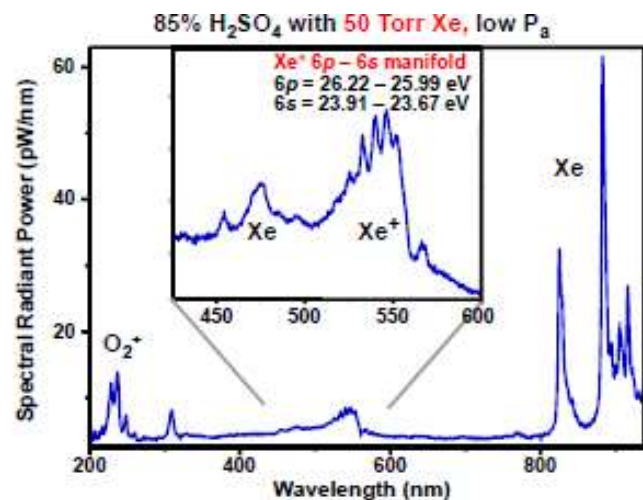
Concentrated H₂SO₄:

exotic system far from applications
but generating beautiful spectra

Very low vapour pressure (from water)
H₂SO₄ 95 wt % : 2.10⁻³ mbar (10,000 times less than water)
⇒ Bubble content = dissolved (rare) gas
⇒ More efficient energy concentration
⇒ Less quenching

SBSL, 85% H₂SO₄, 67 mbar gas

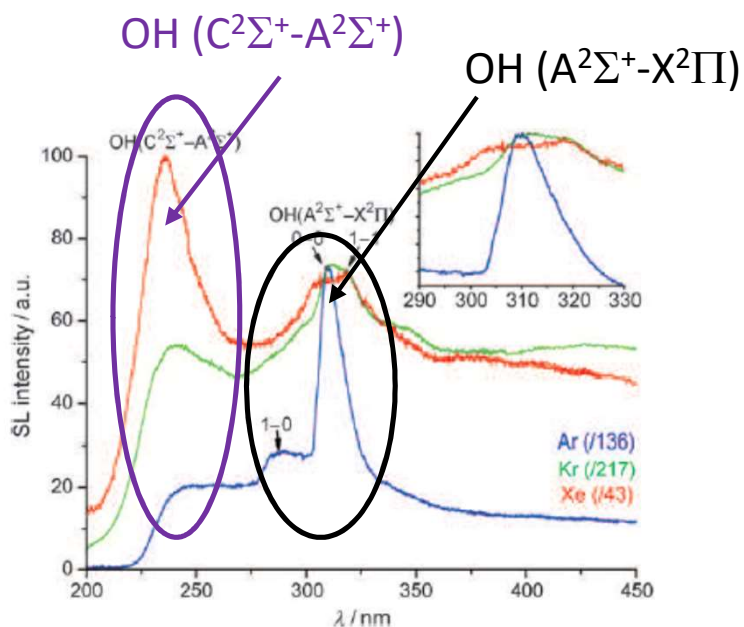
Very high energy levels of emitting species



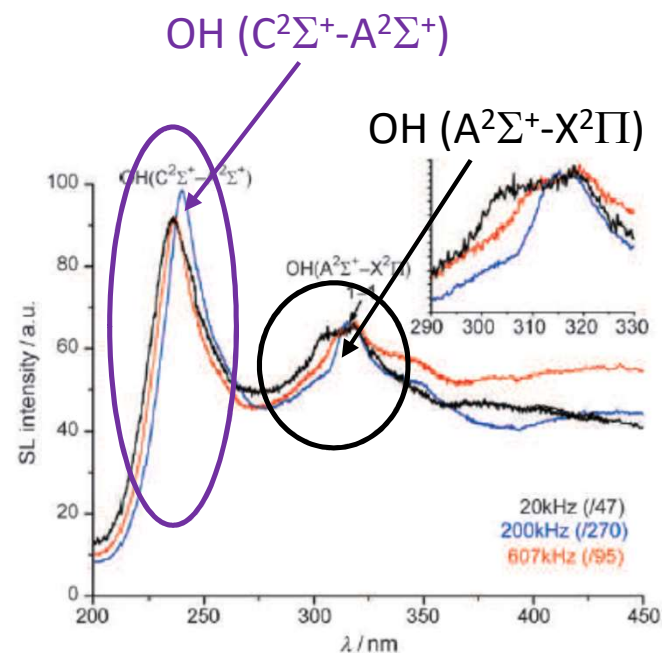
O₂⁺ not thermally generated (in general : electron impact 50-100 eV)

Ion emission extremely rare in SL

In aqueous solutions: no ions
but OH (C-A) transition



MBSL spectra H₂O, Ar/Kr/Xe, 20 kHz



MBSL spectra H₂O, Xe, 20/200/607 kHz

Intensity of OH ($C^2\Sigma^+ - A^2\Sigma^+$): Ar < Kr < Xe and increases with the US frequency

Emission never observed in flames,

only in **discharges** through water vapor (*Shuaibov, Technical Phys. 2009*),

and in **γ -irradiated water** (*Sitharamarao, j. Phys. Chem. 1963*)

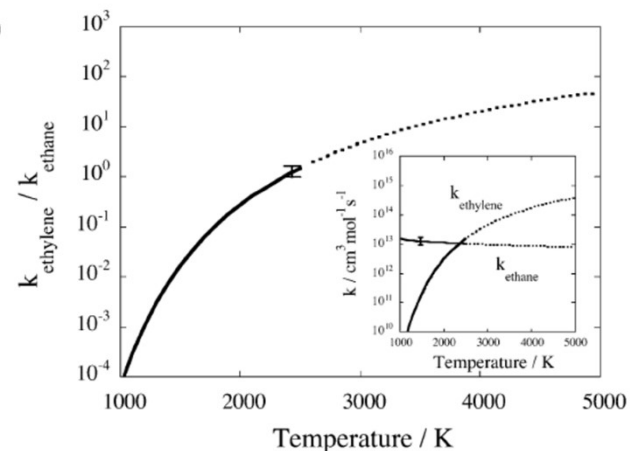
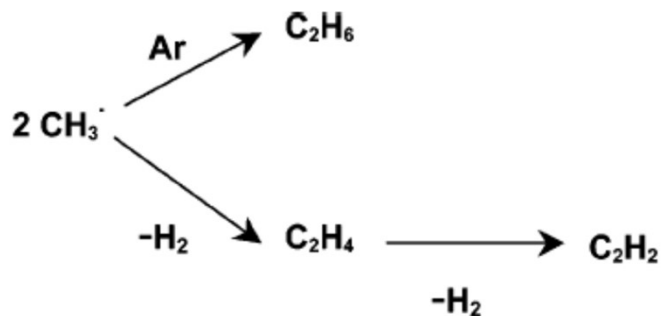
Sonoluminescence spectroscopy in aqueous solutions: what temperatures are estimated?



Mean bubble temperature

Ciawi et al.: Aqueous solutions of *tert*-butanol, Ar

Methyl radical recombination method: temperature estimated from the kinetics of recombination reactions (thermal equilibrium assumed)



Extrapolation of the temperature to zero concentration:

US frequency, kHz	Chemical temperature, K
20	3400 ± 200
355	4300 ± 200
1056	3700 ± 200

Measurement but:

- “*mean chemical temperature*”: spatial and temporal resolutions not well defined
- parallel reaction pathways neglected

Sonoluminescence is emitted at the last stage of collapse
It reflects the most extreme conditions reached

First approach = **temperature derivation from the SL continuum shape**
Approach particularly used in SBSL measurements

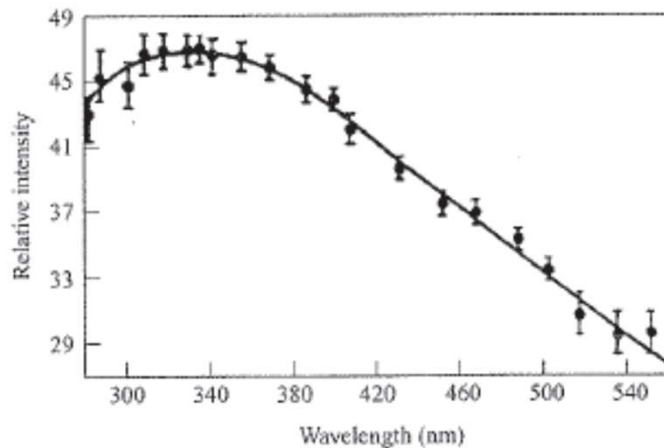


Figure 4.1 Spectral distribution for oxygen-saturated water. Solid line is the theoretical curve for black-body temperature of 8800 K. (Srinivasan and Holroyd (1961).)

Black-body: 8800 K

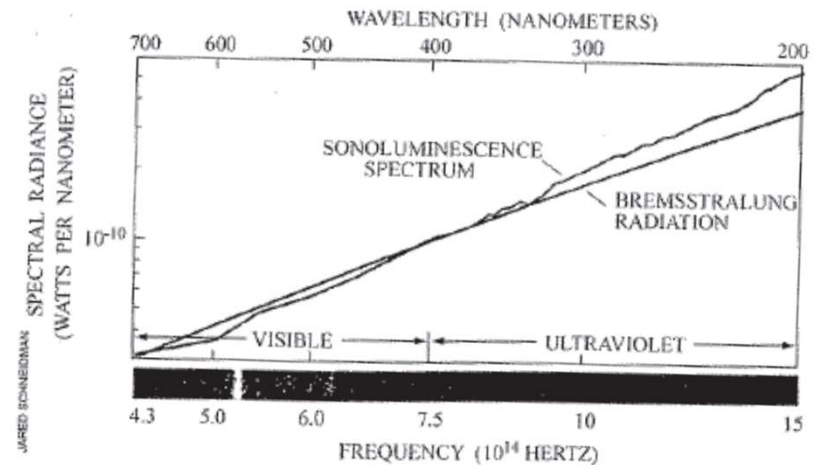


Figure 4.4a Spectrum of sonoluminescence shows that most of the emitted light is ultraviolet. As pointed out by Paul H. Roberts and Cheng-Chin Wu of the University of California at Los Angeles, the signal compares closely with bremsstrahlung radiation – that is, light emitted by a plasma at 100,000 K. (Putterman (1995).)

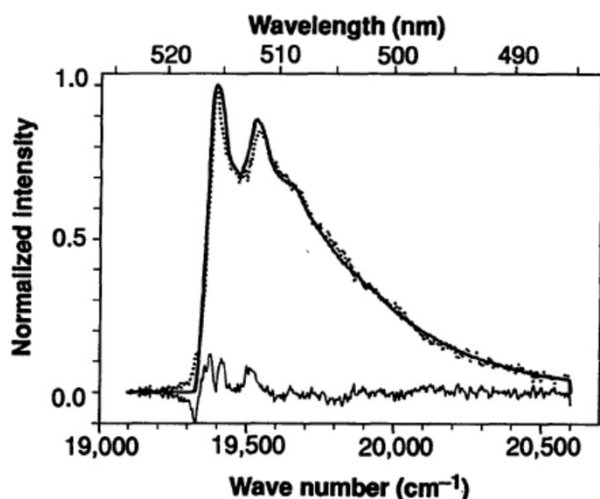
Bremsstrahlung: 100,000 K

Highly model dependent!

And: why should there be only one contribution to the continuum?

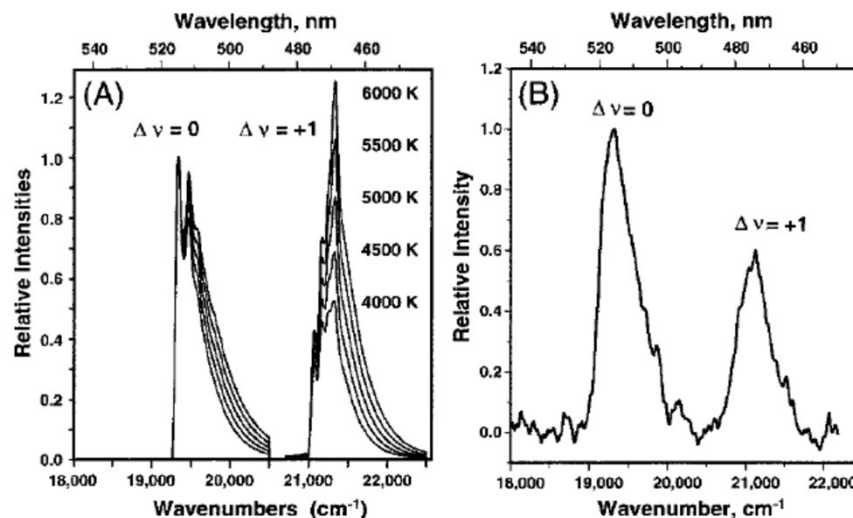
Molecular emissions provide a more direct approach since their shape reflects their excitation level (i.e. the relative populations of excited levels).

The temperature of cavitation,
 Flint & Suslick (1991)
 = the first one
 Silicon oil, Ar, 20 kHz
 C₂ Swan bands



Unicity of temperature was assumed
 5075 ± 156 K

Hot spot conditions during cavitation in water,
 Didenko et al (1999)
 Rare example in aqueous solutions
 Benzene aq. solution, 20 kHz, Ar
 C₂ Swan bands

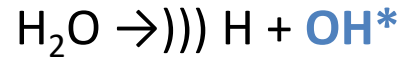


Unicity of temperature was assumed
 4300 ± 200 K

Unicity of the temperature?



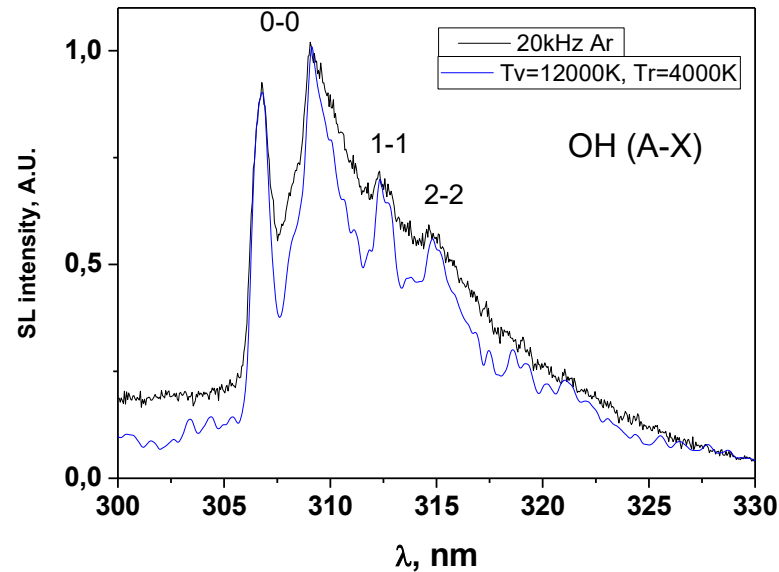
SL spectroscopy of water, Ar, 20 kHz



OH (A-X) emission

Simulations realized with
Specair software

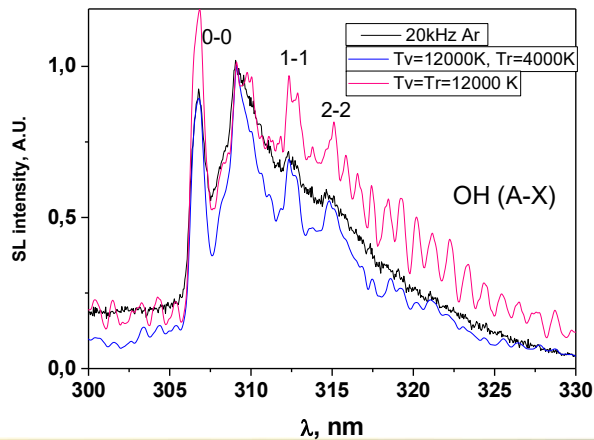
Assuming $p = 500$ bar



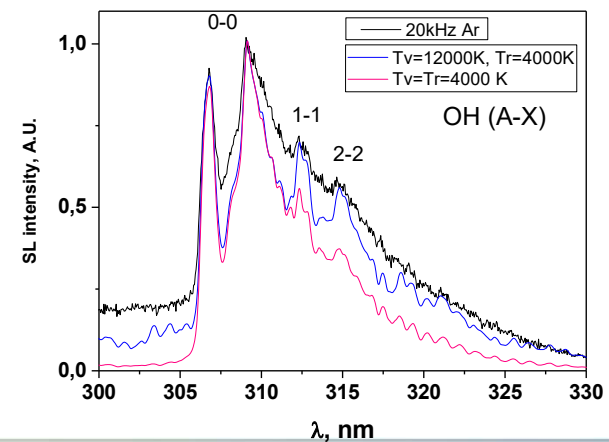
**Need to introduce
different temperatures!**

$T_v = 12000$ K, $T_r = 4000$ K

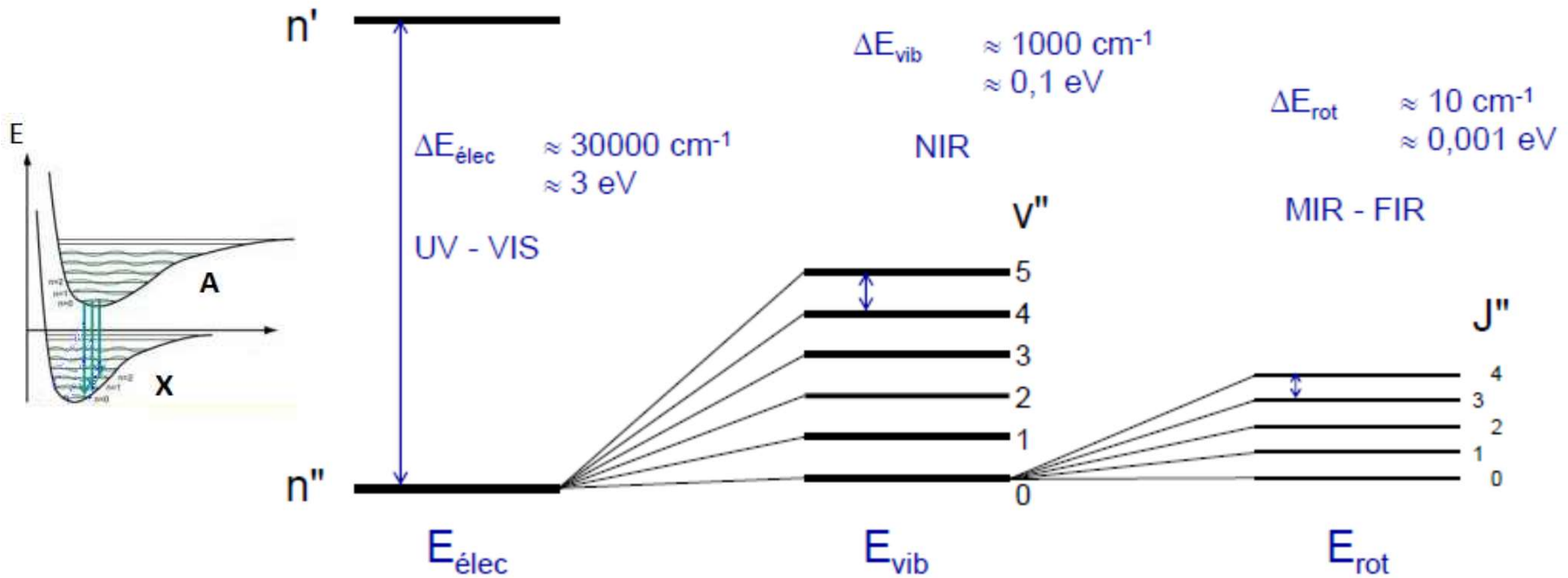
For $T_v = T_r = 12000$ K



For $T_v = T_r = 4000$ K



What do temperatures reflect?



Electronic levels

Relative population of each determined by the electronic temperature T_e

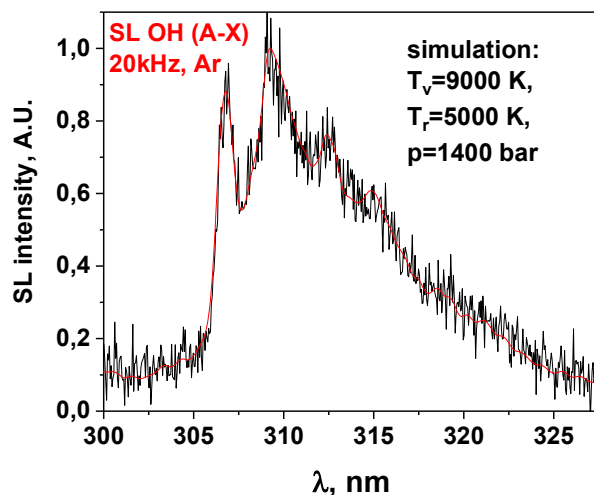
Vibrational levels

Relative population of each determined by the vibrational temperature T_v

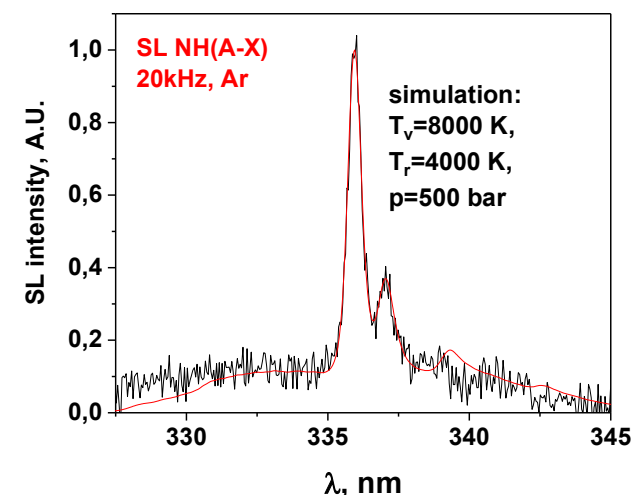
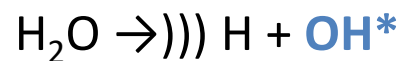
Rotational levels

Relative population of each determined by the rotational temperature T_r

If time scales are too short, thermal equilibrium is not reached
General case in a non-equilibrium plasma: $T_e > T_v > T_r \sim T_{\text{gas}}$



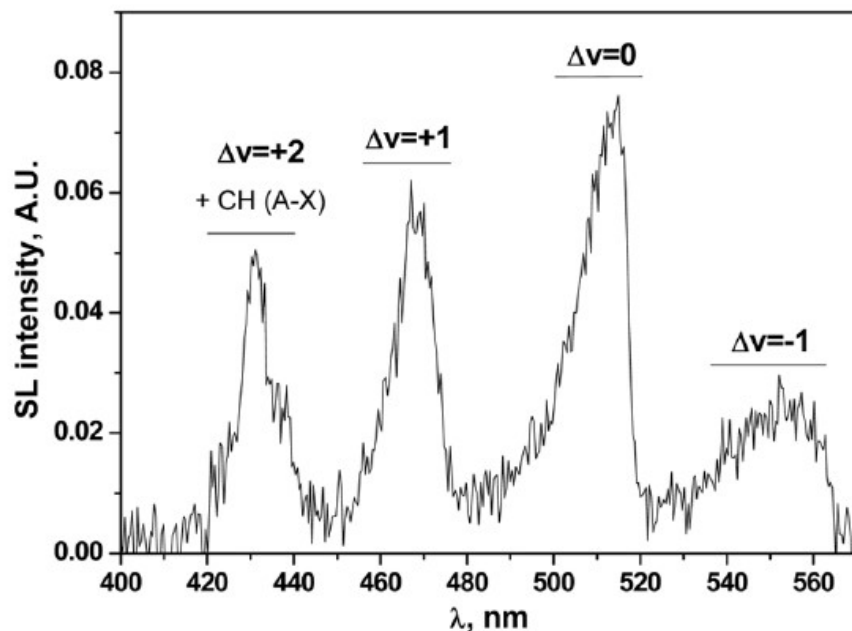
0.1 M $\text{NH}_3 \cdot \text{H}_2\text{O}$, Ar
2 molecular emissions:
OH (A-X) & NH (A-X)



transition	US frequency	$T_v \pm 1000$	$T_r \pm 500$	p , bar
NH ($A^3\Pi-X^3\Sigma^-$)	20 kHz	7000 K	4000 K	500 ± 50
OH ($A^2\Sigma^+-X^2\Pi$)	20 kHz	9000 K	5000 K	1400 ± 200
OH ($A^2\Sigma^+-X^2\Pi$) in water	20 kHz	12000 K	4000 K	1400 ± 200

$T_v > T_r$: non-equilibrium plasma, as in H_2O , Ar
Good agreement between OH & NH temperatures
Lower T_v in ammonia solution compared to water

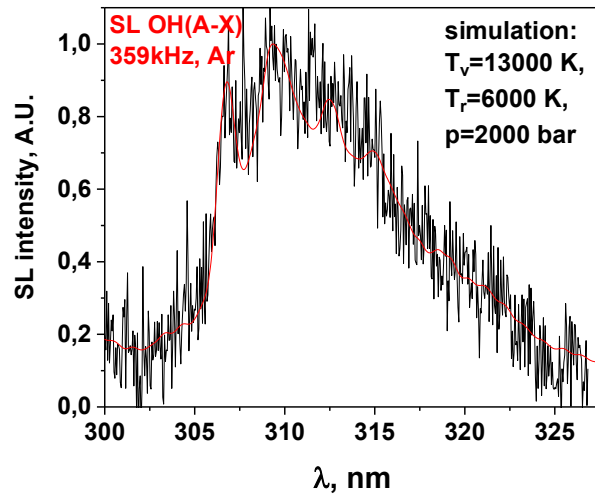
0.1 M t-BuOH, Ar, 20 kHz, 11°C



[t-BuOH] = 0.05 – 0.4 M

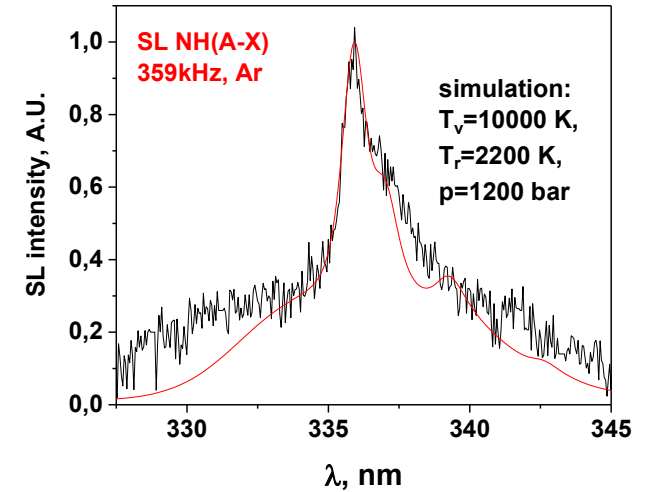
SL spectra present **C₂ Swan bands**
 C₂ (d³Π_g – a³Π_u)

Solution	transition	US frequency	T _v ± 1000	T _r ± 500
water	OH (A ² Σ ⁺ -X ² Π)	20 kHz	12000 K	4000 K
NH ₃ ·H ₂ O	NH (A ³ Π-X ³ Σ ⁻)	20 kHz	7000 K	4000 K
	OH (A ² Σ ⁺ -X ² Π)	20 kHz	9000 K	5000 K
<i>Tert-butanol</i>	C ₂ (d ³ Π _g – a ³ Π _u)	20 kHz	6300 K	4800 K



0.1 M NH₃·H₂O, Ar
359 kHz

OH (A-X) & NH (A-X)
emissions

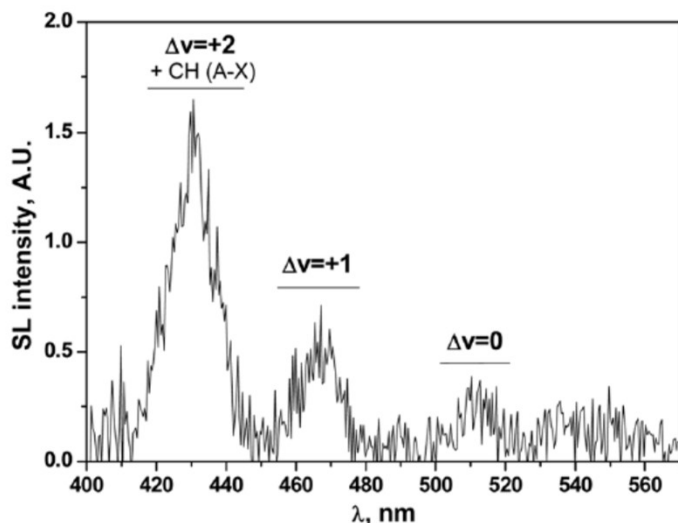


transition	US frequency	T _v , K	T _r , K	p, bar
NH (A ³ Π-X ³ Σ ⁻)	359 kHz	10000 ± 1000	2200 ± 500	1200 ± 100
OH (A ² Σ ⁺ -X ² Π)	359 kHz	13000 ± 2000	6000 ± 1000	2000 ± 500

$$T_v > T_r$$

Good agreement between OH & NH temperatures

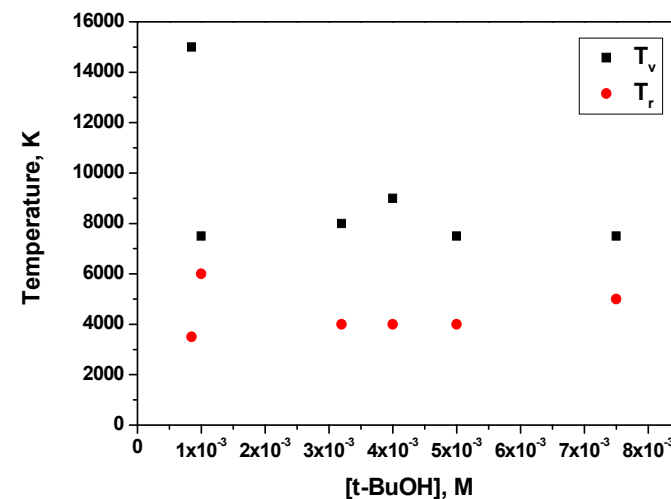
8.5 10^{-4} M t-BuOH, Ar, 362 kHz, 5°C



[t-BuOH] = $1 \cdot 10^{-3}$ – $5 \cdot 10^{-3}$ M; 362 kHz

SL spectra present C₂ Swan bands
C₂ ($d^3\Pi_g - a^3\Pi_u$)

Solution	transition	US frequency	T _v ± 1000	T _r ± 1000
NH ₃ ·H ₂ O	NH (A ³ Π-X ³ Σ ⁻)	359 kHz	10000 K	2200 K
	OH (A ² Σ ⁺ -X ² Π)	359 kHz	13000 K	6000 K
Tert-butanol	C ₂ ($d^3\Pi_g - a^3\Pi_u$)	359 kHz	6300 K	4800 K



The presence of volatile solutes induces a decrease in T_v, in agreement with observations made on ammonia solutions.



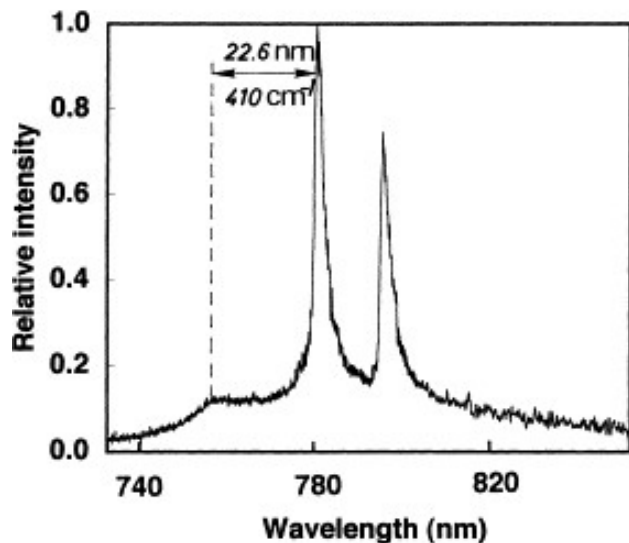
What about the pressure estimation?



Two experimental techniques to determine pressure / intrabubble density:

➤ **Shift in wavelength**

- Lepoint-Mullie (*Ultrason Sonochem* 2001) (Rb*, 20 kHz, Ar) : $d_{rel} = 18 \pm 2$
- McNamara (*Nature* 1999) (Cr*, 20 kHz, Ar, silicone oil) : $d_{rel} = 19 \pm 2$
(corresponding to 300 bar at 4700 K)



		Transitions in Ar gas			
		$^2P_{3/2}-^2S_{1/2}$		$^2P_{1/2}-^2S_{1/2}$	
		(nm)	(cm ⁻¹)	(nm)	(cm ⁻¹)
Shift	H ₂ O	0.41	6.7	0.56	8.9
	1-octanol	0.58	9.5	0.67	10.6
		$\times 10^{-26}$ (m ⁻³) (amagats)		$\times 10^{-26}$ (m ⁻³) (amagats)	
Density	H ₂ O	4.2	15.7	5.0	18.7
	1-octanol	5.4	20.2	5.6	20.9

Lepoint-Mullie (*Ultrason Sonochem* 2001)

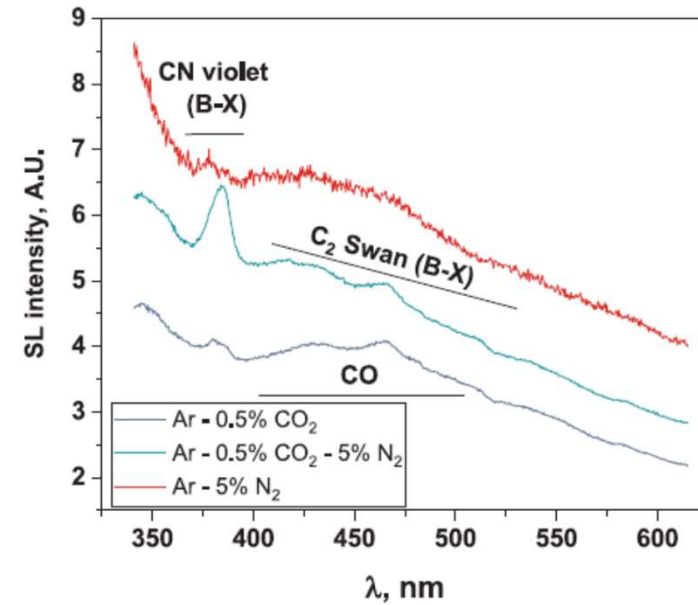
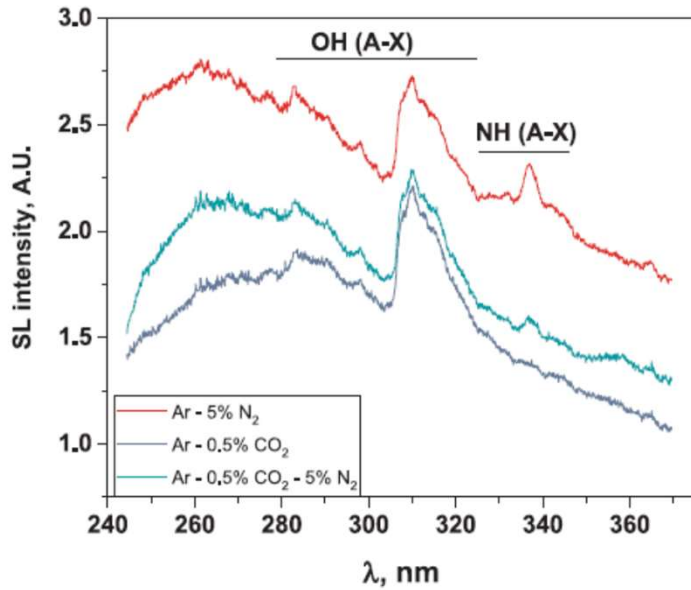
➤ **Peak broadening** (present approach also)

- Hypothesis: pressure = main source of peak broadening
- Sehgal (*J Chem Phys* 1979) (Na*, 460 kHz, Ar): $d_{rel} = 36-50$
- Choi (*J Phys Chem B* 2008) (Na*, 138 kHz, Ar): $d_{rel} = 59,5$ (i.e. 880 bar for 4300 K)
- Kazachek (*Techn Phys Lett* 2009) (Na*, 22/44 kHz, Ar): 800-1200 bar

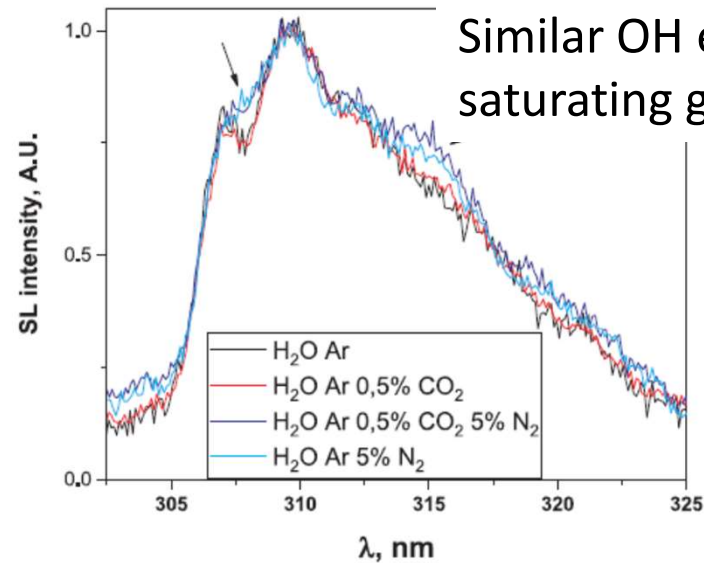
Obtained values are much higher! >> $d_{rel} = 18 \pm 2$ given by Lepoint-Mullie

Other source of peak broadening?

Example of water, 362 kHz Ar mixtures



Observed emissions:
 Ar: OH
 Ar - 5% N₂: OH, NH
 Ar - 0.5% CO₂: OH, CO, (CN)
 Ar - 0.5% CO₂ - 5% N₂: OH, C₂, CN



Similar OH emission for all saturating gases

→ Emissions attributed to NH (c-a) and N₂ (C-B) prevent fitting of OH in some cases

Estimated rovibronic temperatures and effective pressure



Specair software; $T_v, T_r \pm 1000$ K; $P_{\text{eff}} \pm 300$ bar

gas	OH	NH	C ₂	CN
Ar	9,000 K / 6,500 K / 1,000 bar			
Ar – 5% N ₂		10,000 K / 10,000 K / 4,000 bar		
Ar – 0.5% CO ₂ – 5% N ₂			22,000 K / 5,000 K / 6,000 bar	10,000-15,000 K / 10,000 K / 10,000 bar

Large discrepancy in the obtained p_{eff} value depending on the considered species

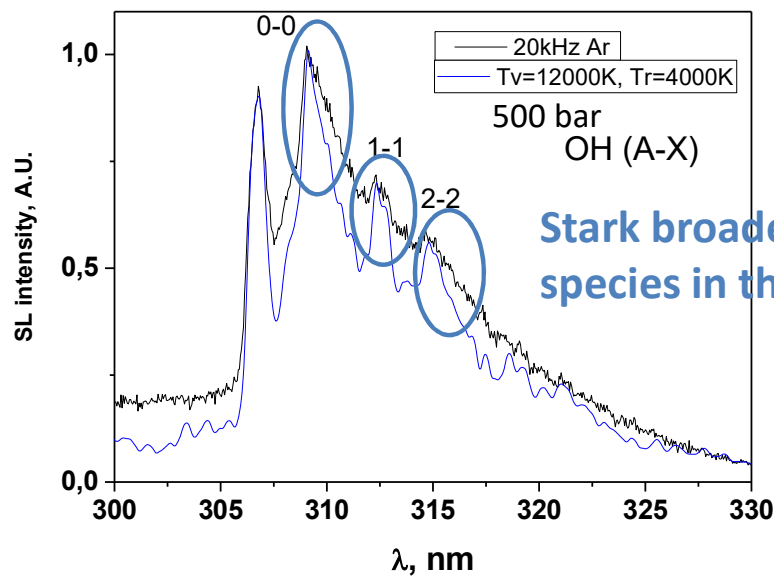
This value of **effective pressure quantifies the broadening** of the emissions.

It is **not a real pressure**.

Another source of peak broadening, the **perturbation by the charged species present in the plasma (Stark effect) must be taken into account**

Neutral species → pressure broadening

Charged species → Stark broadening



Stark broadening: mark of the charged species in the plasma

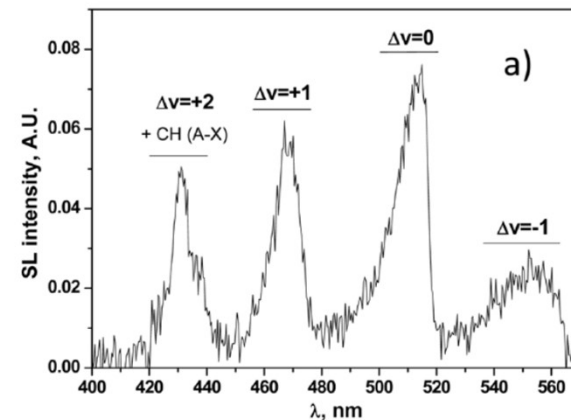
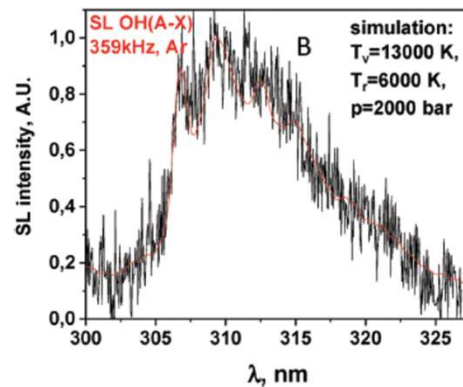
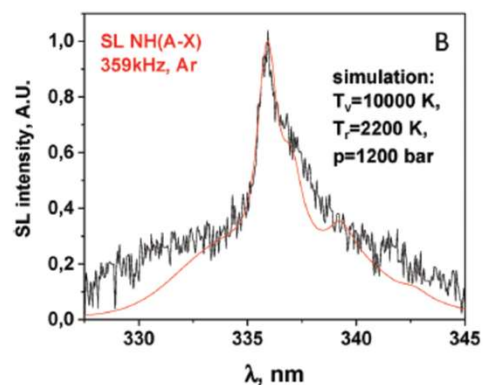
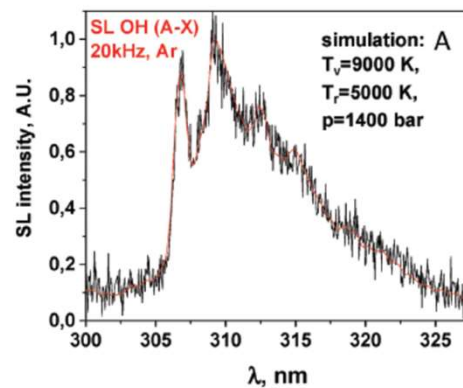
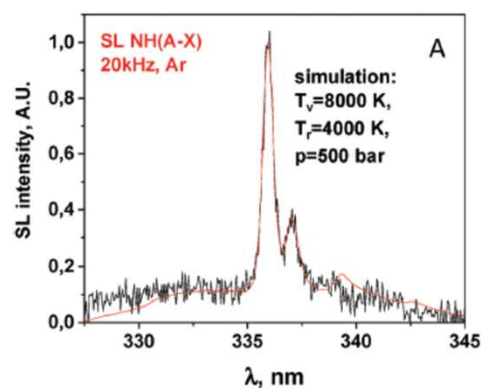
Approach previously adopted in SBSL by Flannigan (*Phys Rev Lett* 2006) (Ar^* , H_2SO_4)

Molecular emissions (MBSL, H_2O) much more complex to model (lack on tabulated data on Stark parameters)

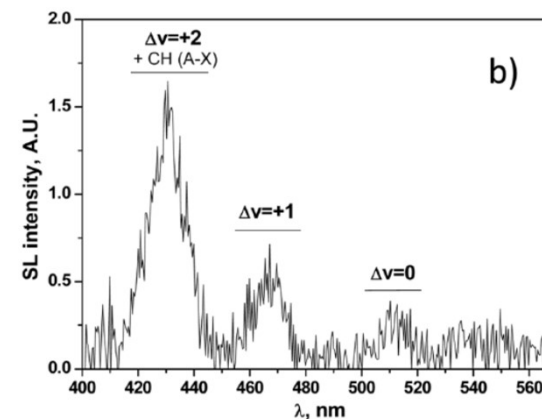
Strong broadening of emissions at HF

Estimated T_v (Specair) HF > BF

More ionized plasma at HF



20 kHz :
6300 /
4800 K



362 kHz :
8000 /
4000 K

t-BuOH aq, 20 kHz 0,1M Ar (a) &
362 kHz 8,5 10⁻⁴M Ar (b)

Fig. 5: Molecular emission of OH ($A^2\Sigma^+-X^2\Pi$) and NH ($A^3\Pi-X^3\Sigma^-$) in SL spectra of 0.1 M ammonia solutions sonicated under Ar at 20 kHz and 359 kHz, and their simulations with Specair



Link SL spectra – sonochemical activity

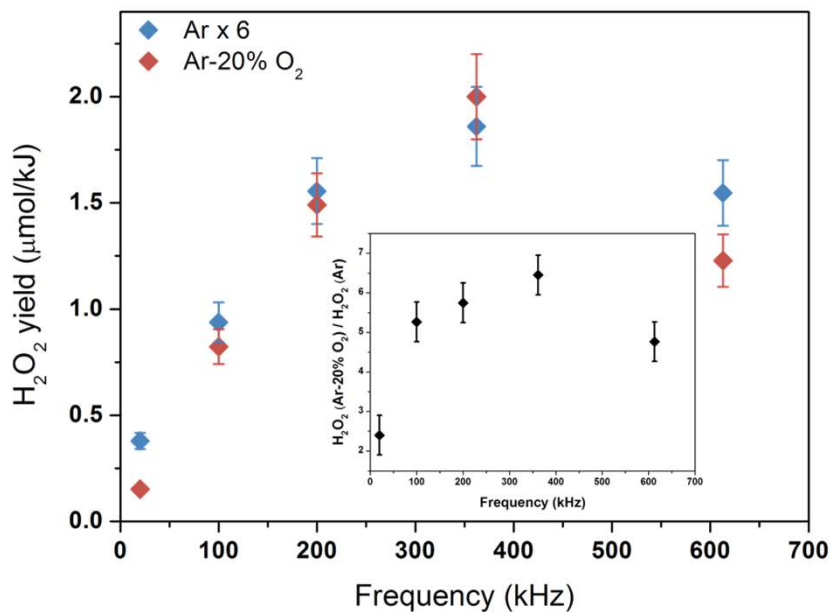


Ar vs. Ar/O₂ at different frequencies

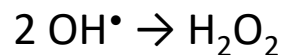
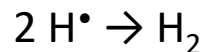
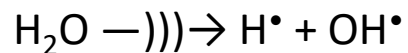


Luminol sonochemiluminescence

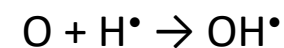
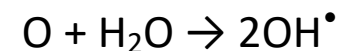
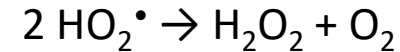
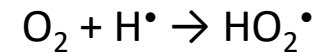
	20 kHz	100 kHz	362 kHz
Ar			
Ar-20%O ₂			



Under Ar:



In the presence of O₂:

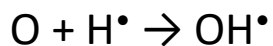
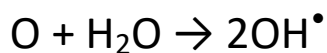
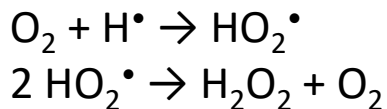


Higher increase in H₂O₂ yield at à HF

Ar vs. Ar/O₂ at different frequencies



Adding 20%O₂ in Ar: LF: H₂O₂ ↑, SL ↓
 HF: H₂O₂ ↑, SL ↑

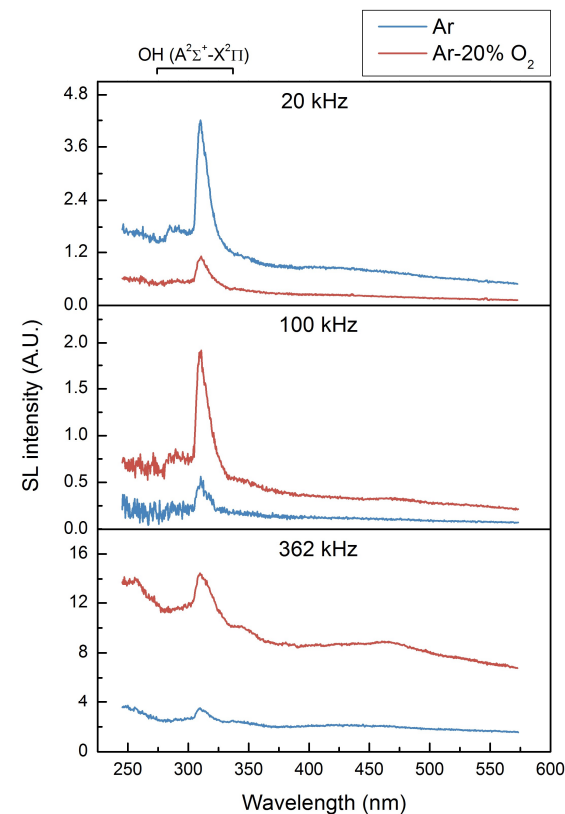
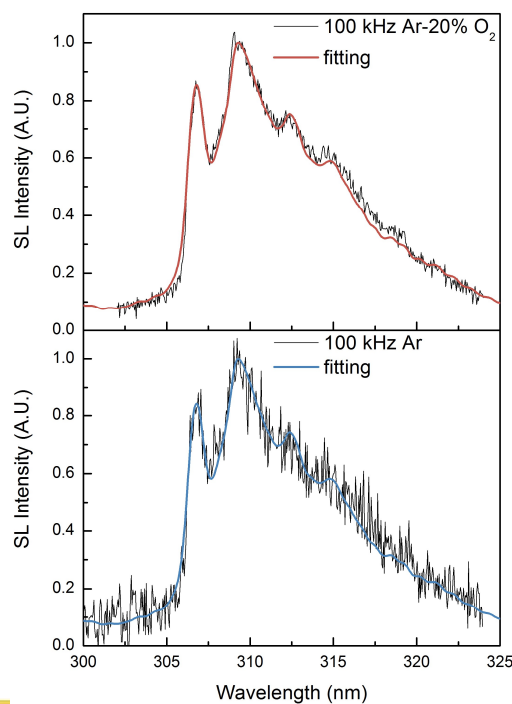
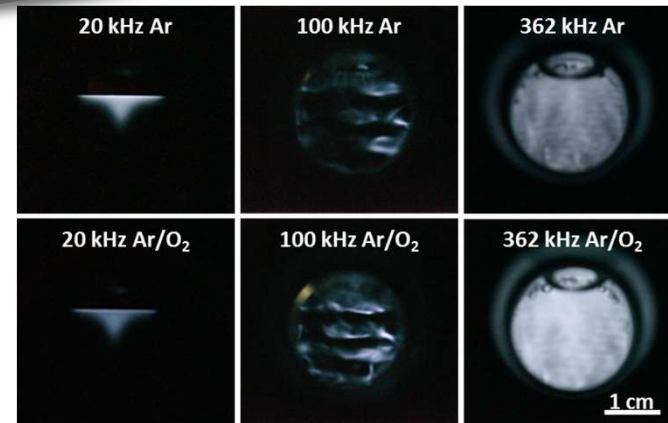
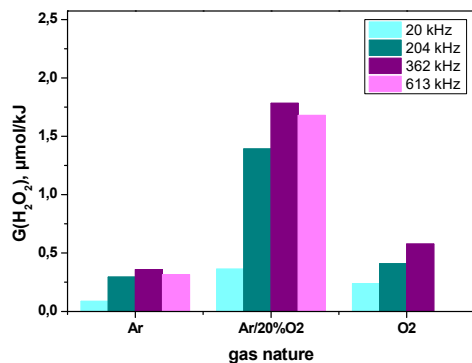


At HF: O₂ dissociation gets important

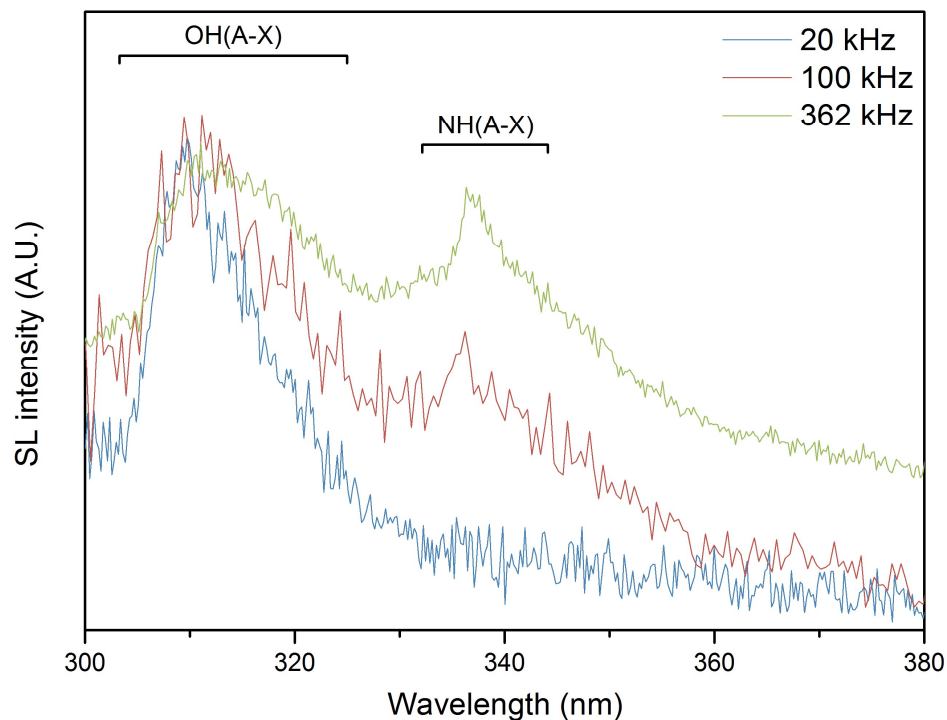
- Different degradation mechanisms can be observed at HF and LF

Different formation pathways for OH result in different T_v

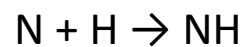
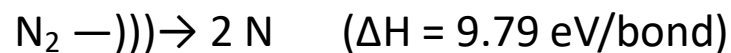
	Ar	Ar/O ₂	
100 kHz	T _v , K	6500 ± 500	7250 ± 250
	T _r , K	4750 ± 250	5200 ± 200



Water Ar-20%N₂
MBSL spectra



20 kHz ($P_{ac} = 33$ W)
 100 kHz ($P_{ac} = 40$ W)
 362 kHz ($P_{ac} = 43$ W)
 18 ° C, 100 mL.min⁻¹ Ar/(20 vol.%)N₂



Dissociation of N₂ at high and intermediate frequencies (≥ 100 kHz)



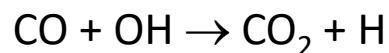
Non-equilibrium sonochemistry?



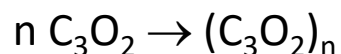
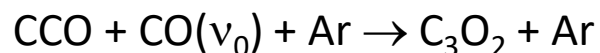
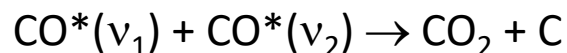
CO disproportionation



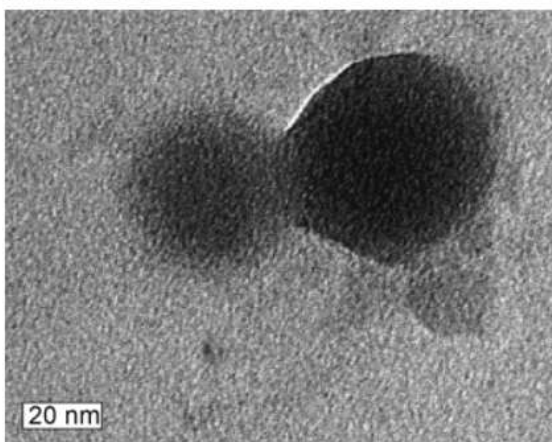
20 kHz sonication of **water saturated with Ar - 20% CO:**



But also formation of a solid similar to $(\text{C}_3\text{O}_2)_n$ polymer



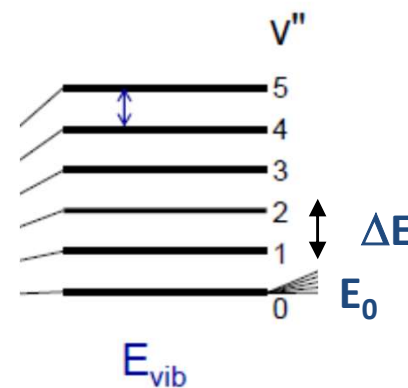
A sound reaction: The ultrasonically driven disproportionation of CO in water causes formation of hydrated polymer of carbon suboxide $(\text{C}_3\text{O}_2)_n$ (see picture) enriched with a heavy ^{13}C isotope. This finding indicates the presence of vibrationally excited CO molecules with a non-equilibrium Treanor distribution inside the cavitating bubble.



$(\text{C}_3\text{O}_2)_n$ solid enriched
in ^{13}C isotope
Non equilibrium

At **equilibrium**, enrichment with **light isotope** expected: it reacts faster due to its higher zero vibrational level energy.

Here the reaction occurs via **vibrational excitation**: heavier isotopes have higher T_v .



20 kHz sonolysis of H₂O/D₂O mixtures saturated with Ar / Xe

Measurement of H₂, D₂, HD in the gas phase (MS)

H/D isotope separation factor α

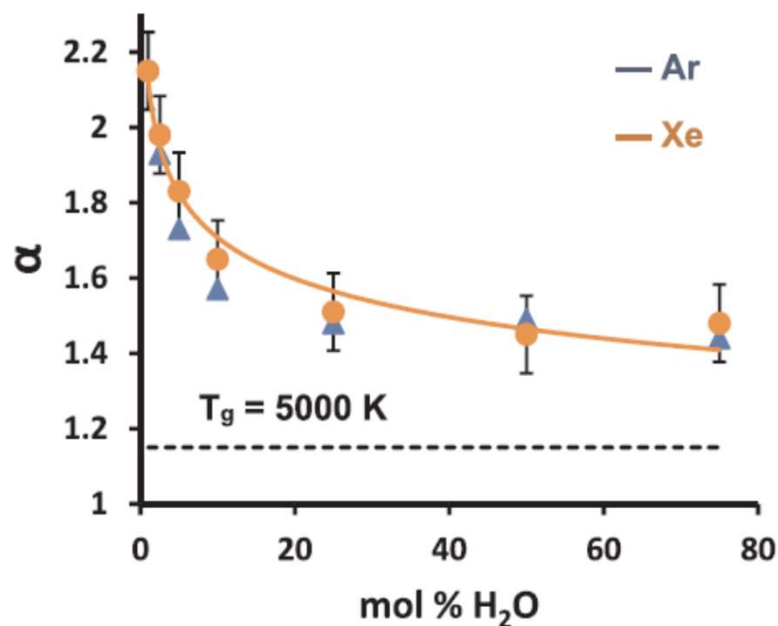


Fig. 3. Dependence of H/D KIE (α) on H₂O concentration in H₂O/D₂O mixtures sparged with Ar (\blacktriangle) and Xe (\bullet). The dotted line shows calculated α value at $T_g = 5000$ K presuming classical behavior.

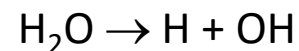
$$\alpha = \frac{\left(\frac{H}{D}\right)_{product}}{\left(\frac{H}{D}\right)_{initial}} = \frac{\left(\frac{I(H_2) + \frac{1}{2}I(HD)}{I(D_2) + \frac{1}{2}I(HD)}\right)}{\left(\frac{H_2O}{D_2O}\right)}$$

Equilibrium case ($T_g = 5000$ K) $\alpha < 1.2$

98% D₂O $\alpha \approx 2$

Enrichment in light isotope

Not only water molecule splitting as usually reported



But more complex plasma mechanism involving e^-_h , D_2O^+ , HD_2O^+ ...

Indoubtedly the **hot spot** formed in cavitation bubbles **is a non-equilibrium plasma**, even in MBSL in aqueous solutions

- **plasma**: presence of **electrons & charged species**
in aqueous solutions, no emission of ions in SL
indirect indication of their presence = strong broadening of emissions (Stark effects)
- **non-equilibrium**: no unique temperature characterizes excited species

Implications of this non-equilibrium in sonochemistry have been illustrated by few experimental studies but would deserve more interest.

Thank you



For your kind attention!